

Stable Difference Schemes with Uneven Mesh Spacings*

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Abstract. We consider a finite-difference approximation to the Cauchy problem for a first-order hyperbolic partial differential equation using different mesh spacings in different portions of the domain. By reformulating our problem as a difference approximation to an initial-boundary value problem, we are able to use the theory of H. O. Kreiss and S. Osher to analyze the L_2 stability of our scheme.

0. Introduction. In solving partial differential equations by finite-difference approximations, there are situations where the solution exhibits large gradients in a localized region. In such problems, one might wish to employ variable mesh patterns. The principal problem is then to obtain consistent difference equations at the interface of the mesh patterns without introducing any type of instability or without loss of overall computational efficiency.

This paper considers the L_2 stability of a finite-difference solution to the Cauchy problem $u_t = Au_x + Bu_y : u(x, y, 0) = \phi(x, y)$, using two different spatial mesh patterns. Problems of this type shall be referred to as *mesh refinement problems* in this paper.

Stability is analyzed by reformulating our problem as a difference approximation to an initial-boundary value problem for a system of partial differential equations. General sufficient conditions for the stability of such systems are due to H. O. Kreiss [3], [4] and S. Osher [6].

Numerical results obtained on New York University's CDC-6600 are presented which yield evidence of stability for cases not analytically treated (e.g. refinements having corners, and refinements using different time steps in different portions of the domain).

I. Formulation of Problem. We shall refer to the following initial-value problem as *Example I*.

$$(1.1) \quad u_t = \alpha u_x : u(x, 0) = \phi(x)$$

for $t > 0$ and all x with α constant and u and ϕ scalar functions.

Consider a numerical solution using two different patterns of mesh spacings. On the right-hand side of the origin, the mesh length is Δx_1 ; on the left-hand side, let the spatial mesh size be Δx_2 . Setting $p = \Delta x_1 / \Delta x_2$, take the mesh on the right-hand side as more *refined* than on the left, i.e. $0 < p < 1$. In the following, let w_i^n be the

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Lax-Wendroff difference approximation [5] to the solution $u(x, t)$ of Example I at $x = j\Delta x_1, t = n\Delta t$. Then, for $j = 1, 2, \dots$, and $n = 0, 1, 2, \dots$,

$$(1.2) \quad w_i^{n+1} = \left(\frac{b_1^2 - b_1}{2}\right)w_{i-1}^n + (1 - b_1^2)w_i^n + \left(\frac{b_1^2 + b_1}{2}\right)w_{i+1}^n, \quad b_1 = \alpha\Delta t/\Delta x_1.$$

To treat this initial-value problem by reduction to an initial-boundary value problem in the quarter plane, it is convenient to introduce a new difference function v_i^n which denotes an approximation to $u(-(j - 1)\Delta x_2, n\Delta t)$ for $j \geq 0$.

Now, with the vector notation

$$(1.3) \quad Z_i^n = \begin{pmatrix} w_i^n \\ v_i^n \end{pmatrix}, \quad j = 0, 1, 2, \dots, \quad n = 0, 1, 2, \dots$$

and with the obvious definition $w_0^n \equiv v_1^n$, the mesh refinement scheme can be expressed as

$$(1.4a) \quad Z_i^{n+1} = \begin{pmatrix} \frac{b_1^2 - b_1}{2} & 0 \\ 0 & \frac{b_2^2 + b_2}{2} \end{pmatrix} Z_{i-1}^n + \begin{pmatrix} 1 - b_1^2 & 0 \\ 0 & 1 - b_2^2 \end{pmatrix} Z_i^n + \begin{pmatrix} \frac{b_1^2 + b_1}{2} & 0 \\ 0 & \frac{b_2^2 - b_2}{2} \end{pmatrix} Z_{i+1}^n$$

for $j = 1, 2, \dots, n = 0, 1, 2, \dots$ (where $b_2 = pb_1$). Actually, there remains to prescribe an approximation to $u(0, t)$ at each time step, since a formula for v_0^n is still lacking. To derive an interface condition, we make the natural requirement that our difference scheme be of second-order accuracy at the point $x = 0$. Using the consistency conditions derived from Taylor's theorem, one could obtain directly a three-point asymmetrical scheme up to second-order accuracy. However, we prefer an equivalent formulation which expresses v_0^n by second-order extrapolation in terms of w_1^n, v_1^n and v_2^n . We complete the formulation of the mesh refinement problem as a matrix difference approximation to an initial-boundary value problem by expressing this interface condition in the matrix form,

$$(1.4b) \quad Z_0^n = \begin{pmatrix} w_0^n \\ v_0^n \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ \alpha_1 & \alpha_2 \end{pmatrix} \begin{pmatrix} w_1^n \\ v_1^n \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 0 & \alpha_3 \end{pmatrix} \begin{pmatrix} w_2^n \\ v_2^n \end{pmatrix}, \quad n = 0, 1, 2, \dots,$$

where

$$\alpha_1 = \frac{2}{p(1 + p)}, \quad \alpha_2 = \frac{2(p - 1)}{p}, \quad \alpha_3 = \frac{1 - p}{1 + p}.$$

Because our difference equations were obtained by reflection of the left-hand side, (1.4) is consistent with the hyperbolic system

$$\begin{pmatrix} w \\ v \end{pmatrix}_t = \begin{pmatrix} \alpha & 0 \\ 0 & -\alpha p \end{pmatrix} \begin{pmatrix} w \\ v \end{pmatrix}_x$$

with compatibility conditions $w(0, t) = v(0, t)$. We shall refer to the system of difference equations (1.4a, b) as *System I*.

For initial-boundary value problems of the type formulated above, Kreiss [3],

[4] and Osher [6] have determined algebraic conditions which when satisfied guarantee stability of our difference approximation.

In the algebraic verification of stability, for most of our mesh refinement problems, the following immediate corollary of the Main Theorem and Lemma 3 of Kreiss [4] has been used repeatedly. It is assumed that the reader is familiar with the terminology of that paper.

COROLLARY I. *The approximation \mathcal{G} is stable if*

1. *Assumptions 1-4 (of [4]) are satisfied, and if*

2. *for all $z, |z| \geq 1, \text{Det } E(\tau_i(z)) = 0$ implies that some $l + 1$ characteristic roots have absolute value equal to one.*

Note, $E(\tau_1(z))$ refers to the system (1.12) for $|z| \geq 1$ and for $z = 1$ to (1.12') both of [4]. Also note that, for our System I, $l = 1$.

II. Stable Mesh Refinement Schemes. In this section, we verify stability for the mesh refinement problem System I and introduce generally applicable interface conditions which yield stable mesh refinement schemes.

THEOREM I. *System I is a stable approximation for the Cauchy problem of Example I for $0 < |b_1| < 1$.*

Proof. For $|b_1| < 1$, assumptions 1-4 are satisfied since the Lax-Wendroff scheme is a stable, dissipative scheme. To check stability, find all eigenvalues λ , of the mesh refinement problem, by expressing the eigensolutions as

$$(2.1a) \quad z_i^n = \lambda^n g_i; \quad \sum_{i=1}^{\infty} |g_i|^2 < \infty; \quad \lambda \neq 1, \quad |\lambda| \geq 1,$$

where $g_i = g_i(\lambda)$ is determined using the general solution of the characteristic equations

$$(2.2) \quad \lambda = \left(\frac{b_i^2 + (-1)^i b_i}{2} \right) \frac{1}{\tau_i} + (1 - b_i^2) + \left(\frac{b_i^2 + (-1)^{i+1} b_i}{2} \right) \tau_i, \quad i = 1, 2.$$

The g_i take the form

$$(2.1b) \quad g_i(\lambda) = g_i = \begin{pmatrix} \beta_1 & \tau_1^i \\ \beta_2 & \tau_2^i \end{pmatrix}, \quad \beta_i \text{ constants,}$$

since each τ_i (the characteristic roots) are separated by the unit circle for $|\lambda| \geq 1, \lambda \neq 1$. The β_i here are referred to as $\sigma = (\sigma_1 \sigma_2)'$ in [4]. We now show that all nontrivial solutions of the form (2.1), when substituted into the boundary condition (1.4b), imply that $\tau_1 = \tau_2 = 1$. Thus, stability follows by Corollary I.

Substitution of (2.1) into the boundary condition yields

$$(2.3) \quad (a) \quad \beta_1 = \beta_2 \tau_2, \quad (b) \quad \beta_2 = \alpha_1 \beta_1 \tau_1 + \alpha_2 \beta_2 \tau_2 + \alpha_3 \beta_2 \tau_2^2,$$

for a nontrivial solution, $\beta_2 \neq 0$. Eliminating β_1 and β_2 gives

$$(2.4) \quad p(1 - p)\tau_2^2 + 2\tau_1\tau_2 + 2(p^2 - 1)\tau_2 - p(1 + p) = 0.$$

Equating the expressions for λ from (2.2), yields, upon clearing fractions,

$$(2.5) \quad 2b_1^2(1 - p^2)\tau_1\tau_2 + \tau_1[p^2b_1^2 + pb_1 + (p^2b_1^2 - pb_1)\tau_2^2 - (b_1^2 + b_1)\tau_1\tau_2] - (b_1^2 - b_1)\tau_2 = 0.$$

Solve for $\tau_1\tau_2$ and τ_1 from (2.4) and substitute into the above to find, after some rearrangement,

$$\left(\frac{b_1^2 - b_1}{2}\right)p^2(p^2 - 1)(\tau_2 - 1)^4 = 0.$$

For a reasonable problem, this can only be satisfied if $\tau_2 = 1$. Now, setting $\tau_2 = 1$ in (2.4) yields that $\tau_1 = 1$. Q.E.D.

We now present two different interface conditions whose stability is easily demonstrated and which are useful in higher dimensions. Consider the formula for extrapolation of a sufficiently smooth function $v(x)$,

$$(2.6) \quad v_0 = \sum_{i=1}^m (-1)^{i+1} \binom{m}{i} v_i,$$

where $v_j = v(j\Delta x)$ (see e.g. [2]). Alter the format of System I so that $v_i^n \sim u(-j\Delta x_2, n\Delta t)$, for $j = 0, 1, 2, \dots$; and so $v_0^n = w_0^n$ for all n . Using (2.6) to define v_0^n , refer to the resulting mesh refinement approximation scheme as System II.

THEOREM II. *System II is a stable approximation to Example I when $\alpha < 0$. If $\alpha > 0$, then replacing v_i^n in (2.6) by w_i^n results in a stable mesh refinement scheme for $|b_1| \leq 1$.*

Proof. The algebraic verification of stability now follows trivially in the obvious matrix formulation.

An interface condition, which has the advantage of being stable regardless of the sign of α , and is also suitable for the case of higher dimensions, can be had by taking $p = \Delta x_1/\Delta x_2 = 1/M$ where M is an integer. Assuming a definition for v_i^n and w_i^n as in System I, we see that v_0^n coincides with w_M^n for $n = 0, 1, \dots$. This suggests use of the coarse mesh difference equation at the interface. This boundary condition in matrix form becomes

$$(2.7) \quad \begin{pmatrix} w_0^n \\ v_0^n \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} w_1^n \\ v_1^n \end{pmatrix} + \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} w_M^n \\ v_M^n \end{pmatrix}, \quad n = 0, 1, 2, \dots$$

We shall refer to the difference approximation formed by using (2.7) in conjunction with the difference equations (1.4a) as System III.

THEOREM III. *System III is a stable approximation to Example I for $|b_1| \leq 1$.*

Proof. Again, because of the separation of the roots, to check stability by Kreiss' theorem, try a solution of the form (2.1). Substitution into the boundary condition gives for nontrivial solutions

$$\beta_1 = \beta_2\tau_2 \quad \text{and} \quad \beta_2 = \beta_1\tau_1^M.$$

A nontrivial solution exists only if β_1 and β_2 are both different from zero. Eliminating the β_i 's, we find

$$(2.8) \quad \tau_2\tau_1^M = 1.$$

For a bounded solution, the above implies that $|\tau_2| = |\tau_1| = 1$. Thus, stability follows from Corollary I.

We remark that the above formulation can be used to show that the interface conditions of Systems II and III can be used to match any 3-point dissipative difference scheme (even different ones on each side) along an interface [1] and retain stability.

III. Equations with Two Space Dimensions. Consider the analogous mesh refinement problem for

$$(3.1) \quad u_t = au_x + bu_y, \quad a, b \text{ constants, for all } x \text{ and } y \text{ and } t \geq 0.$$

Let the plane $x = 0$ be the interface between two mesh patterns. In the right half-plane, take mesh spacings in the x and y directions to be Δx_1 and Δy_1 , respectively. In the left half-plane, denote these respective mesh spacings by Δx_2 and Δy_2 . We restrict our attention to the case where $p_1 = \Delta x_1/\Delta x_2 = 1/N$ and $p_2 = \Delta y_1/\Delta y_2 = 1/M$, where N and M are positive integers.

Let $W_{i,k}^n$ and $V_{i,k}^n$ denote approximations to $u(j\Delta x_1, k\Delta y_1, n\Delta t)$ and

$$u(-(j - 1)\Delta x_2, k\Delta y_2, n\Delta t),$$

respectively, for $n = 0, 1, 2, \dots, j = 0, 1, 2, \dots, k = 0, \pm 1, \pm 2, \dots$ (see Fig. 1). As in the one-dimensional problem, one needs to express $W_{0,k}, V_{0,k}$ as a linear combination of the values at neighboring points. From the definitions of our mesh functions,

$$(3.2) \quad V_{0,k} = W_{N,kM}, \quad W_{0,kM} = V_{1,k}, \quad k = 0, \pm 1, \pm 2, \dots$$

At each time step, we still have to prescribe $W_{0,kM+i}^n$ for $k = 0, \pm 1, \pm 2, \dots, i = 1, 2, \dots, M - 1$. In general, one could prescribe such points as a linear combination of the $V_{1,i}$ values, say,

$$(3.3) \quad W_{0,kM+i}^n = \alpha_{i,-t} V_{1,k-t}^n + \dots + \alpha_{i,0} V_{1,k}^n + \dots + \alpha_{i,s} V_{1,k+s}^n$$

for s, t natural numbers.

Now, consider a mesh refinement problem for a 9-point dissipative approximation to (3.1), (e.g. the Lax-Wendroff scheme [5]). Then

$$(3.4) \quad \begin{aligned} (a) \quad V_{i,k}^{n+1} &= \sum_{\alpha=-1}^1 \sum_{\beta=-1}^1 C_{\alpha,\beta}^{(1)} V_{i+\alpha,k+\beta}^n \\ (b) \quad W_{i,k}^{n+1} &= \sum_{\alpha=-1}^1 \sum_{\beta=-1}^1 C_{\alpha,\beta}^{(2)} W_{i+\alpha,k+\beta}^n \end{aligned}$$

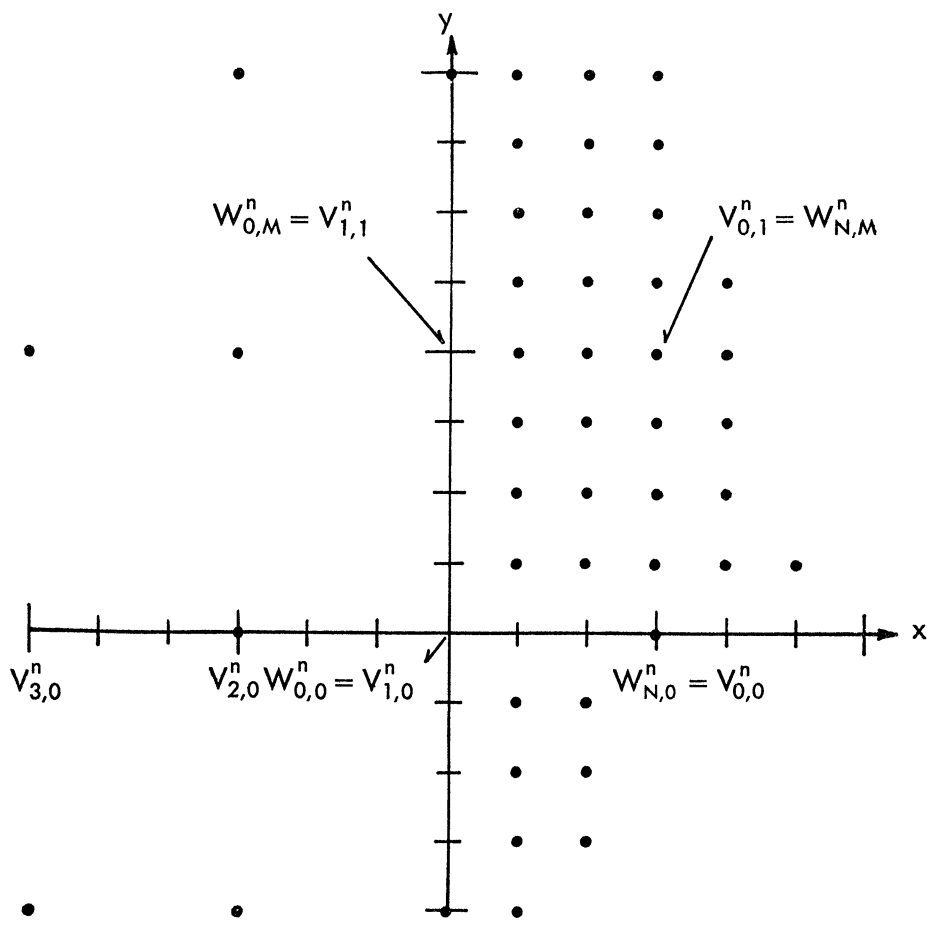
$$n = 0, 1, 2, \dots, j = 1, 2, 3, \dots, k = 0, \pm 1, \pm 2, \dots,$$

where the $C_{\alpha,\beta}^{(i)}$ are functions of $\Delta t, \Delta x_i, \Delta y_i$. Then, (3.2)–(3.4) is a well-defined approximation to (3.1). Again, we treated the mesh refinement problem as a system of difference equations for an initial-boundary value problem. Because of the more complicated nature of the boundary condition in two dimensions, we found it necessary to introduce a vector of dimension $M + 1$ corresponding to the data $(V_{i,k}^n, W_{i,kM}^n, \dots, W_{i,kM+M-1}^n)$. Then, reformulating our difference equations consistent with a system of $M + 1$ partial differential equations, we are able to prove [1].

THEOREM IV. *If the difference schemes (3.4) are stable for the pure Cauchy problem, then the mesh refinement scheme (3.2)–(3.4) is a stable approximation.*

Proof. Because of the diagonal structure of our systems, it suffices to study all eigensolutions of the form

$$Z_{i,k}^n = \lambda^n y_0^k g_i, \quad |y_0| = 1, \quad \sum_{j=1}^{\infty} |g_j|^2 < \infty,$$



$$N = 3, M = 4$$

FIGURE 1

and show that Corollary I is satisfied for all $y_0, |y_0| = 1$. One finds that the conditions (3.2), when expressed in the indicated matrix form, lead directly to Eq. (2.8). Thus, arguing as in Theorem III, stability is demonstrated.

We note that following Kreiss [3] and [1], the matrix analogues of Theorems II-IV hold for Lipschitz continuous diagonal systems.

IV. Computational Results. Calculations were performed using the mesh refinement schemes of Theorems I-IV. The primary purpose was to compare the relative accuracy of the different interface conditions discussed above. We tested our methods on problems where the analytical solution was available.

a. *One-Dimensional Periodic Case.* Table 1 lists the results of several calculations

using the Lax-Wendroff scheme for the initial-value problem of Example I with $\alpha \equiv 1$ and solution $u(x, t) = \sin 4\pi(x + t)$. We used the periodicity of our solution to restrict our calculations to the interval $[0, 1]$ with refined region $I_R \equiv \{x: \frac{1}{3} \leq x \leq \frac{2}{3}\}$. By $e(v)$ and $e(w)$, we denote the maximum error of our scheme from the solution in the unrefined region and the interior of I_R , respectively.

To describe if any additional error due to our unusual interface conditions propagated into I_R , we compute a quantity defined by

$$e_{100} \equiv \max_{j=2,3,\dots,100} \{|u(x, t) - w_j^n|: t = n\Delta t, x = \frac{1}{3} + (j - 1)\Delta x_1\},$$

since the error from $x \geq \frac{2}{3}$ cannot reach into more than 400 grid points of I_R in 400 time steps.

In the mesh refinement problem, a small time step is needed for stability in I_R . However, using the same time step in the unrefined region will reduce accuracy there. To rectify this situation we also considered using different time steps in the different regions. The only difficulty in using this approach is at the interface. The interfaces for Systems I and III were treated by first advancing $w_1^{n/p+i}$ and $w_{501}^{n/p+i}$, for $i = 1, 2, \dots, 10$. Then, the remaining is advanced in the usual manner until all points are at the same time level. Rows 7 and 8 of Table 1 show the improved results, and a comparison of rows 1-3 (or 7, 8) with row 6 shows that the interface conditions used at $x = \frac{1}{3}$ do not propagate much error into I_R since the e_{100} agree to 7 places in all entries.

In Table 2 we present the results of computations performed for Example I with

$$(4.1) \quad \begin{aligned} \alpha = \alpha(x) &= \frac{1}{2}, & 0 \leq x \leq \frac{1}{3}, \\ &= \frac{2 + \sin 6\pi x}{4}, & x \in I_R, \\ &= \frac{1}{2}, & \frac{2}{3} \leq x \leq 1. \end{aligned}$$

The quantities displayed in Table 2 are defined exactly in the same way as in Table 1.

TABLE 1

$$\alpha \equiv 1$$

$$\Delta t = 1/1750; \quad \Delta t' = 1/175; \quad \Delta x_2 = 1/150; \quad \Delta x_1 = 1/1500$$

	$t, \text{ time}$	$e(v)$	$e(w)$	e_{100}
1) System I	$400 \Delta t$	3.332×10^{-3}	1.591×10^{-3}	6.3×10^{-6}
2) System II ($m = 3$)	$400 \Delta t$	3.332×10^{-3}	1.656×10^{-3}	6.3×10^{-6}
3) System III	$400 \Delta t$	3.332×10^{-3}	1.629×10^{-3}	6.7×10^{-6}
4) No Refinement (Δx_2)	$400 \Delta t$	3.333×10^{-3}		
5) No Refinement (Δx_2)	$40 \Delta t'$	8.909×10^{-4}		
6) No Refinement (Δx_1)	$400 \Delta t$	8.9×10^{-6}	8.9×10^{-6}	6.3×10^{-6}
7) Uneven Time (I)	$40 \Delta t'$	8.909×10^{-4}	4.493×10^{-4}	6.3×10^{-6}
8) Uneven Time (III)	$40 \Delta t'$	8.909×10^{-4}	4.493×10^{-4}	6.3×10^{-6}

TABLE 2

$\alpha = \alpha(x)$ as in (4.1)
 $\Delta t = 1/1250; \Delta t' = 1/125; \Delta x_2 = 1/150; \Delta x_1 = 1/1500$

	$t, \text{ time}$	$e(v)$	$e(w)$	e_{100}
1) System I	400 Δt	7.136×10^{-3}	1.397×10^{-3}	1.000×10^{-3}
2) System II ($m = 3$)	400 Δt	7.134×10^{-3}	2.764×10^{-3}	1.000×10^{-3}
3) System III	400 Δt	9.993×10^{-3}	1.401×10^{-3}	1.000×10^{-3}
4) No Refinement (Δx_2)	400 Δt	1.180×10^{-2}		
5) No Refinement (Δx_2)	40 $\Delta t'$	2.358×10^{-2}		
6) No Refinement (Δx_2) $\Delta t^* = 15/1250$	27 Δt^*	2.901×10^{-2}		
7) Uneven Time (I)	40 $\Delta t'$ = 400 Δt	7.668×10^{-3}	1.652×10^{-3}	1.000×10^{-3}
8) Uneven Time (III)	40 $\Delta t'$ = 400 Δt	7.668×10^{-3}	1.611×10^{-3}	1.000×10^{-3}
9) Uneven Time (III)	27 Δt^* = 405 Δt	4.841×10^{-3}	1.553×10^{-3}	1.086×10^{-3}

b. *Two-Dimensional Periodic Case.* Calculations were performed for the initial-value problem $u_t = -u_x + u_y$ with periodic solution $u(x, t) = \sin 2\pi(x - t) \cdot \cos 2\pi(y + t)$ restricted to the unit square $0 \leq x, y \leq 1$. The refined region was taken as $D_R = \{(x, y): \frac{1}{3} \leq x, y \leq \frac{2}{3}\}$. As in Table 1, we describe the accuracy with the analogous quantities $e(v)$ and $e(w)$, the maximum errors of our scheme in the unrefined region and D_R , respectively. To describe the accuracy of the refined scheme away from the interfaces, we compute the error in the interior of D_R , e_{int} , as the maximum error in the interior square $\{(x, y): \frac{4}{5} \leq x, y \leq \frac{6}{5}\}$. Table 3 lists the results of our calculations using the Lax-Wendroff scheme with our mesh refinement scheme

TABLE 3

$\Delta t = 1/1750; \Delta x_2 = \Delta y_2 = 1/45; \Delta x_1 = \Delta y_1 = 1/225$

	$t, \text{ time}$	$e(v)$	$e(w)$	e_{int}
1) Ordinary refinement	100 Δt	1.182×10^{-3}	1.143×10^{-3}	4.36×10^{-5}
2) No refinement ($\Delta x_2, \Delta y_2$)				
a) $\Delta t = 1/1750$	100 Δt	1.165×10^{-3}		
b) $\Delta t' = 5/1750$	20 $\Delta t'$	1.104×10^{-3}		
c) $\Delta t^* = 1/175$	10 Δt^*	8.895×10^{-4}		
3) Uneven time	20 $\Delta t'$ = 100 Δt	1.104×10^{-3}	1.033×10^{-3}	4.36×10^{-5}
4) Uneven time	10 Δt^* = 100 Δt	8.895×10^{-4}	7.822×10^{-4}	4.36×10^{-5}

using centered third-order polynomial interpolation along the interface. This is compared with ordinary Lax-Wendroff with three choices of time steps. We then compute with these time steps, using uneven time steps in a manner similar to the one-dimensional case. Our computations show that the interface condition does not introduce any additional errors, and the error in the unrefined region compares with the error for ordinary no refinement.

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